

# Method for parametric imaging of attenuation by intravascular optical coherence tomography

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**Abstract:** Catheter-based intravascular optical coherence tomography (IVOCT) is a powerful imaging modality for visualization of atherosclerosis with high resolution. Quantitative characterization of various tissue types by attenuation coefficient (AC) extraction has been proven to be a potentially significant application of OCT attenuation imaging. However, existing methods for AC extraction from OCT suffer from the challenge of variability in complex tissue types in IVOCT pullback data such as healthy vessel wall, mixed atherosclerotic plaques, plaques with a single component and stent struts, etc. This challenge leads to the ineffectiveness in the tissue differentiation by AC representation based on single scattering model of OCT signal. In this paper, we propose a novel method based on multiple scattering model for parametric imaging of optical attenuation by AC retrieval from IVOCT images conventionally acquired during cardiac catheterization. The OCT signal characterized by the AC is physically modeled by Monte Carlo simulation. Then, the pixel-wise AC retrieval is achieved by iteratively minimizing an error function regarding the modeled and measured backscattered signal. This method provides a general scheme for AC extraction from IVOCT without the premise of complete attenuation of the incident light through the imaging depths. Results of computer-simulated and clinical images demonstrate that the method can avoid overestimation at the end of the depth profile in comparison with the approaches based on the depth-resolved (DR) model. The energy error depth and structural similarity are improved by about 30% and 10% respectively compared with DR. It provides a useful way to differentiate and characterize arterial tissue types in IVOCT images.

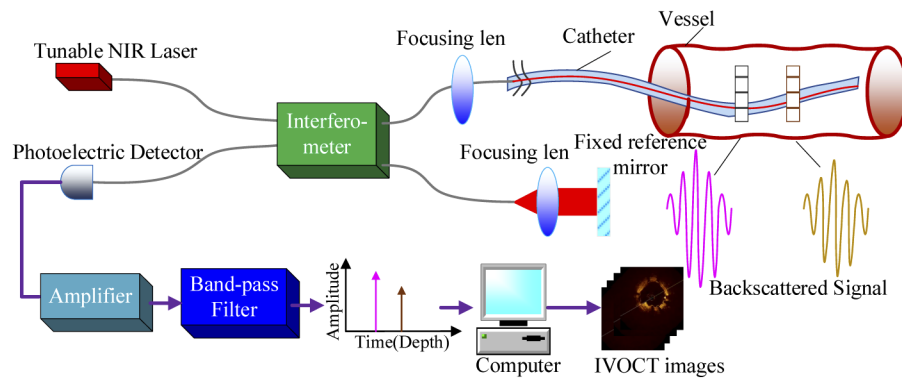
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## 1. Introduction

### 1.1. Backgrounds

Intravascular optical coherence tomography (IVOCT) plays an important role in the diagnosis of coronary arterial atherosclerosis and the guidance of stent implantation. It is a catheter-based imaging modality capturing real-time pullback sequences of transverse and volumetric scanning images of arterial vessels with high resolution [1]. It provides detailed information regarding morphological structures and microstructures of vessel wall and atherosclerotic plaques based on backscattering of the incident near-infrared (NIR) laser in a spectral range of 1250-1350 nm [2]. Figure 1 shows the physical overview of a spectral-domain IVOCT system.

Interpretation of tomographic structural images allows for qualification of dissimilar tissue compositions. However, a comparison mechanism of physiological information in addition to morphological structures is required to obtain parameters with clinical diagnostic value that aid the differentiation of tissue components along with determination of plaque vulnerability. Quantification of tissue optical properties has been proven to be able to assist OCT image interpretation. The attenuation coefficient (AC) has been shown a significant indicator of optical



**Fig. 1.** Schematic diagram of a spectral-domain IVOCT system.

properties describing the extinction of the detected OCT signal with the depth due to optical absorption and scattering [3]. The ACs extracted from IVOCT B-scans or three-dimensional (3-D) pullback data can be used to differentiate thin fibrous cap, calcification and lipid rich plaque and characterize the formation of neointima after stent implantation [3,4]. The combination of OCT attenuation with texture features and geometric structures segmented from structural images has been demonstrated to improve the accuracy of tissue characterization and vulnerable plaque identification [5,6].

### 1.2. Related work

Based on the fact that the backscattered light is exponentially attenuated with the detection depth in tissues, the AC is usually estimated from the measured backscattered signal by curve fitting (CF) through single A-line analysis [7–9] or 3-D analysis [10–12]. The estimation can be optimized by extending the fitting window in those regions with homogeneously distributed optical properties [13]. Moreover, CF has been combined with machine learning to improve the accuracy of tissue characterization [14,15]. The CF methods assume a constant attenuation along a part of the A-scan. Hence, they are most suitable for use on the tissue layers with well-defined layer structure, which subsequently reduces the automation of the methods. Besides, for thin layer structures, AC extraction by CF is unstable and sensitive to the noise in the measured signal. The small variations in the ACs cannot be accurately identified, limiting the OCT potential of high resolution.

In recent years, the depth-resolved (DR) technique has been an efficient tool for AC extraction from OCT with pixel-level resolution and high automation. Vermeer *et al.* [16] designed a DR method for the first time which assumes complete attenuation of all incident light in the range of the imaging depth, the absence of absorption and a constant scattering anisotropy. However, overestimation is observed in the areas where the backscattered light is not fully attenuated, such as eyelid area or hair fiber. Beside this, Vermeer's method takes the maximum detection depth as the upper limit of the integral of the backscattered signal, leading to errors near the end of an A-line due to incorrect background subtraction. Furthermore, it assumes a constant confocal function, resulting in inaccurate estimates when the focal plane is within the sample or on the sample surface.

Amaral *et al.* [17] extended Vermeer's DR model to a general model by quantifying the attenuation of the NIR light with the transmittance. Their results show that the ACs of various tissue types are estimated by inputting an appropriate transmittance, even for those thin tissue layers with the thickness of 100  $\mu\text{m}$ . However, this method does not take into account the multiple scattered light, leading to wrong estimates from the weak OCT signals at the tissue boundaries.

Liu *et al.* [18] developed a DR-CF approach by combining the conventional DR model with CF to determine the cutoff depth. It preserves the CF's advantage of high tolerance to noise and avoids wrong estimates caused by the weak backscattered signals at the tissue boundaries. However, it does not take into account the confocal effect and requires the focal plane to be located above the sample, reducing the signal-to-noise ratio (SNR) of the measurements and making some certain applications impractical.

Smith *et al.* [19] designed a depth-resolved confocal (DRC) method which allows for accurate estimation of the ACs when the focusing plane is located within the sample by taking into account both the confocal function and the sensitivity fall-off. But, overestimation is still observed in the presence of incomplete attenuation of the incident light in deep tissues. To address Vermeer's assumption of full attenuation, Liu *et al.* [20] developed an optimized depth-resolved estimation (ODRE) method to implement AC extraction in any depth range regardless of whether the light is completely attenuated or not. ODRE is inapplicable to arterial vessels in IVOCT images because OCT signals attenuate severely in arterial tissues with large ACs such as lipid plaques, leading to an inaccurate threshold of the SNR.

Both DRC and ORDE require the prior knowledge of the confocal function parameters of the OCT setup employed (i.e., the focal plane position and apparent Rayleigh range), which makes some applications impractical. To address this issue, Dwork *et al.* [21] proposed a fully automated DRC method called autoConfocal method. It obtains the confocal function parameters automatically from the differences in the features between two images of the same sample captured at slightly different view angles. However, increasing errors are observed when the focal plane is far away from the sample because of the nearly constant confocal function. Besides, this method is unsuitable for use on IVOCT because it is infeasible to image the same vascular cross section from different views by using a single catheter pullback.

For clarity, Table 1 presents the overview on the major methods of OCT-based AC extraction.

### 1.3. Objective of The Work

In this paper, a novel model-based method is proposed for parametric imaging of attenuation by extracting ACs in arterial cross sections from transverse IVOCT scanning images. The given data is the conventionally acquired IVOCT gray scale images satisfying the DICOM digitalization standard as well as the configuration parameters provided by the imaging system including the incident light power, central wavelength and full width at half maximum (FWHM). Pixel-wise AC estimates are achieved by iteratively minimizing an error function regarding the measured and modeled backscattered signal which is the output of a forward model based on Monte Carlo (MC) simulation. This method provides a general scheme for AC estimates without the premise of complete attenuation of the incident light in tissues, as opposed to the DR-based approaches. Experiments of numerical simulations and clinical images are conducted to verify its validity. The influence of initial plans of iteration is discussed. The experiment of comparing with the DRC method is also conducted.

Table 1. Overview of Major Methods of OCT-Based AC Extraction

Methods		Assumptions		Advantages	Limitations		Applications		
CF	Single A-line analysis [7–9]	Constant AC along a part of an A-line		Simplicity; high tolerance to noise	Low automation; inapplicable to very thin tissue layers or layers close to the surface; a large amount of data required for accurate fitting		Homogeneous tissue regions and layers (except very thin layers) with well-defined layer structure		
	3D CF [10–12]	Homogeneous optical properties of tissues within the fitting window							
	CF with extended fitting window [13]								
DR	CF combined with machine learning [14,15]	Complete attenuation of all incident light through the imaging depth	Constant scattering anisotropy; absence of absorption; a backscattering coefficient linearly varying with AC; a constant confocal function; a focal plane above the sample	Pixel-wise estimation; high automation	Overestimation in the areas where the backscattered light is not fully attenuated; errors near the end of an A-line; inaccurate estimates in the case of a focal plane within the sample or on the sample surface		Homogeneous and heterogeneous multi-layered tissues except thin-layer tissues		
			A known appropriate transmittance	Applicable to thin tissue layers	Wrong estimates from weak signals at tissue boundaries			All tissue types with known transmittance	
			A focal plane above the sample	Avoid wrong estimate caused by weak backscattered signals at tissue boundaries; high tolerance to noise	Reduced SNR of measurements; inapplicability in some certain applications				Multi-layered tissues
	Improved DR based on transmittance [17]	Absence of optical absorption	A reflectivity proportional to AC	Allow the focusing plane to be located within the sample	Require prior knowledge of confocal parameters (i.e., focal plane position and apparent Rayleigh range)	Overestimation in the case of incomplete attenuation of incident light in deep tissues	Multi-layered tissues; inapplicable to coronary arterial vessels in IVOCT images		
	DR-CF [18]				Require sufficient homogeneity in the bottom tissue layer; require a suitable region identified at the bottom of the image			Increasing errors when the focal plane is far away from the sample	
	DRC [19]								
	ORDE [20]	Focal plane is located close to or within the sample		Allow for AC estimation in any depth range regardless of full attenuation	Automatic estimate of confocal parameters from OCT imagery; allow for pixelwise AC extraction without prior knowledge of OCT system parameters				
	autoConfocal DR [21]								



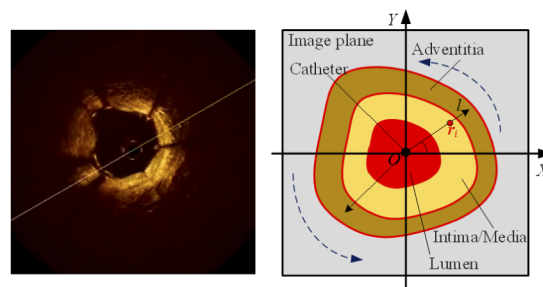
## 2. Method

Our method aims to find the AC distribution in a vascular cross section from an IVOCT gray scale image by solving the forward problem iteratively, updating the ACs at each iteration, until the output of the solver matches the measured OCT signal in a framework of least-squares minimization. That is, the squared error between a measured OCT signal and the output of a forward model is minimized by adjusting the AC. This section gives the detailed steps.

### 2.1. Forward modeling of IVOCT signal

AC extraction from OCT requires a quantitative model to describe the OCT signal. Two general models have been proposed, i.e., single-scattering (SS) model [3,4] and multiple-scattering (MS) model [22,23]. Model selection depends on the tissue type of interest. The SS model assumes that the incident light propagates forward in the medium and is attenuated by absorption and scattering satisfying the Lambert-Beer's law until it is reflected back to the source by backscattering. It is based on the premise of single backscattering of photons, thus it is suitable for the AC extraction of a weakly scattering medium or the superficial thin layer of the highly scattering tissue. The MS model considers multiple backscattering events and large detection depth. Till now, there have been three approaches taking MS into account in OCT. In the order of increasing complexity and number of parameters, they are stochastic approach based on MC simulation [24–26], analytic approach based on extended Huygens–Fresnel model [27] and full electromagnetic wave modeling based on Maxwell's equations [28]. In comparison with the MS model, the SS model has a key advantage of simplicity [3], hence it is the most commonly used in AC extraction such as CF and DR-based methods. However, clinical observations have demonstrated that multiple scattering can be expected in highly forward scattering tissues such as blood and atherosclerotic plaques [3,29].

In this work, we employ MC simulation to construct the forward model of OCT signal. The model includes three parts: Gaussian beam generation, photon propagation in the medium and backscattered signal simulation. For simplicity, we ignore the interface reflection between the optical elements as well as the phase change during light propagation [30]. As illustrated in Fig. 2, in an imaging plane, the catheter remains in the image center, although it does not stay fixed in the lumen center along the lumen length. The lumen, vessel wall intima/media, plaques and adventitia are surrounding the catheter center from inside to outside. A rectangular coordinate system  $XOY$  and a polar coordinate system  $\theta$ - $l$  are established in the imaging plane where the origin  $O$  is located at the image center.



**Fig. 2.** A typical IVOCT B-scan image and schematic diagram of a vascular cross section.

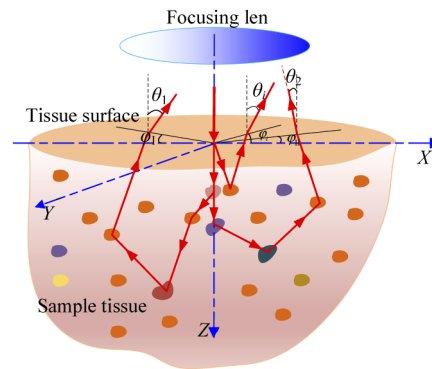
#### 2.1.1. Gaussian beam generation

The position and direction of the photon emission in the Gaussian beam is simulated by MC method. Then, the propagation trajectory of the photon in the focusing system is traced generating

the distribution function on the back surface of the lens. More details on the principle and code of MC simulation of Gaussian beam generation can be found in Ref. [31].

### 2.1.2. Photon propagation in the tissues

On a microscopic scale, each vessel wall layer presents uniform optical property, while their microstructures are regarded as heterogeneous turbid medium. MC simulation of photon propagation depicts the propagation trajectory of photons in a weakly absorbing and strongly scattering media by regarding photons as typical particles and ignoring the properties of polarization and phase of light. As illustrated in Fig. 3, photons are launched perpendicular to the tissue surface. After multiple scattering and absorption, partial photons escape from the medium and are finally received by the detector.



**Fig. 3.** Schematic diagram of scattering events of photons in tissue.  $\theta$  and  $\varphi$  denote the deflection angle and azimuthal angle, respectively.

As illustrated in Fig. 4, the detailed steps are as follows.

### (1) Photon initialization

A single photon of weight  $W$  is launched vertically in the medium surface at an initial position  $(0,0,0)$ . Its direction cosine and weight are initialized as  $(0,0,1)$  and 1, respectively.

## (2) Sampling of photon step-size

Assume the free path of a photon between two successive elastic scattering events is determined by the Poisson probability density function with respect to the AC and the photon free path [30]. The free path-length, i.e., the step-size of photon propagation is given by [31]

$$s = -(\ln \xi)/\mu_t, \quad (1)$$

where  $\xi$  is a computer-generated random number uniformly distributed in the interval  $[0,1]$  and  $\mu_t$  is the AC describing the decay of the incident light due to the tissue optical properties. The pathlength  $L_i$  of the  $i$ -th photon is obtained by the product of its total steps and the refractive index of the medium.

### (3) Photon energy absorption

The partial energy is absorbed by the medium once a photon moves by the distance  $s$  to the scattering site, resulting in the reduction in the photon weight  $W$ , i.e.,  $\Delta W = \mu_a W / \mu_t$ , where  $\mu_a$  is the optical absorption coefficient [33].

#### (4) Photon scattering

After absorption the photon undergoes the scattering event causing it is redirected. The azimuthal angle  $\varphi$  is uniformly distributed from 0 to  $2\pi$  and is obtained by direct sampling, i.e.,  $\varphi=2\pi\xi$ . The cosine of the deflection angle  $\theta \in [0, \pi]$  follows the Henyey-Greenstein probability distribution function [32] as follow,

$$p(\cos \theta) = \frac{1 - g^2}{2(1 + g^2 - 2g \cos \theta)^{3/2}}. \quad (2)$$

Direct sampling gives [33]

$$\cos \theta = \begin{cases} \frac{1}{2g} \left[ 1 + g^2 - \left( \frac{1-g^2}{1-g+2g\xi} \right)^2 \right], & g > 0 \\ 2\xi - 1, & g = 0 \end{cases}, \quad (3)$$

where  $g$  is the scattering anisotropy factor ( $g=1$  for multiple scattering in multi-layered vessel wall tissues implying that all the light is scattered in the forward direction [22,34]).

The penetration depth of the photon in the medium is given by

$$r = \frac{1}{2} \sum_j^n s_j \mu_{z,j}, \quad (4)$$

where  $n$  denotes the total times of a single photon scattering in the medium,  $s_j$  denotes the step size of the  $j$ -th photon and  $\mu_{z,j}$  is the direction cosine of the  $j$ -th scattering,

$$\mu_{z,j} = -\sin \theta \cos \varphi \sqrt{1 - \mu_{z,j-1}^2} + \mu_{z,j-1} \cos \theta. \quad (5)$$

#### (5) Ending of photon propagation

Tracking of the photon ends when its weight is less than the threshold ( $10^{-6}$ ) or when it escapes the geometry boundaries of the medium.

#### 2.1.3. OCT signal simulation

The MC simulation of photon propagation outputs the weights and pathlengths of the photons collected from the medium surface. Then, the OCT signal backscattered by tissues in terms of the radial depth  $r_i$  is determined by [30]

$$I_{\text{OCT}}(r_i) = I_0 \sum_{i=1}^{N_{\text{ph}}} \sqrt{W_i} \cos \left[ \frac{2\pi}{\lambda_c} (2r_i - L_i) \right] \exp \left[ -\left( \frac{2r_i - L_i}{l_{\text{coh}}} \right)^2 \right], \quad (6)$$

where  $l_{\text{coh}}$  is the coherence length of the probe radiation,

$$l_{\text{coh}} = (2 \ln 2 / \pi) \cdot (\lambda_c^2 / \Delta \lambda) \quad (7)$$

Here,  $I_0$  is a constant determined by the instrumental properties of the OCT system,  $N_{\text{ph}}$  is the total number of photons launched (typically  $\sim 10^7$ ),  $W_i$  is the weight of  $i$ -th detected photon with optical pathlength  $L_i$ ,  $2r_i$  is the optical pathlength in the reference arm,  $\lambda_c$  is the central wavelength of the incident radiation, and  $\Delta \lambda$  is the FWHM of the laser spectrum. Eq. (6) accounts for the speckle structure caused by the overlap of interference patterns. Speckles are eliminated in simulations by using the following expression [30] instead of Eq. (6),

$$I_{\text{OCT}}(r_i) = I_0 \sum_{i=1}^{N_{\text{ph}}} \sqrt{W_i} \exp \left[ -\left( \frac{2r_i - L_i}{l_{\text{coh}}} \right)^2 \right] \quad (8)$$

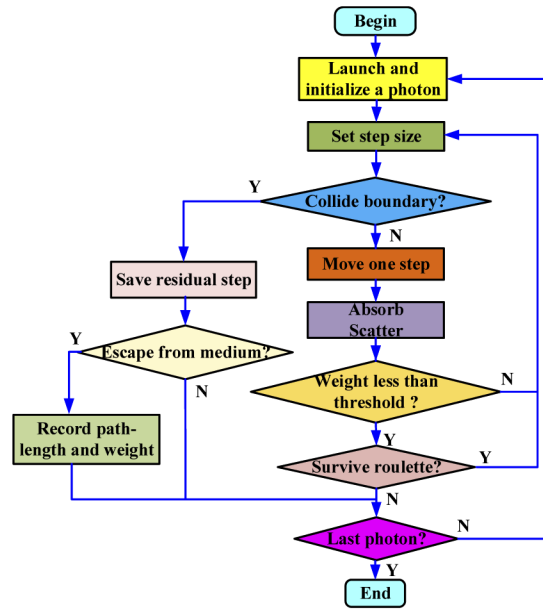


Fig. 4. Flow chart of MC simulation of photon propagation in a medium.

## 2.2. AC estimation by iterative optimization

The standard output of a commercially available IVOCT system is sequential pull-back gray scale images with 8-bit gray levels overlaid with golden color maps as shown in Fig. 2. The intensity of the measured backscattered signal at each radial depth in a vascular cross section is directly proportional to the gray value of the corresponding image pixel,

$$I_m(r_i) = g(r_i)/255, \quad (9)$$

where  $I_m(r_i)$  is the measured OCT signal and  $g(r_i)$  is the 8-bit gray value (ranging from 0 to 255) of the pixel at  $r_i$  in a standard IVOCT B-scan image.

An error function regarding the measured and modeled OCT signal is defined as follow,

$$f(\mu_t(r_i)) = I_{\text{OCT}}(r_i) - I_m(r_i), \quad (10)$$

where  $f(\cdot)$  is the error function to be minimized,  $\mu_t(r_i)$  is the AC and  $I_{\text{OCT}}(r_i)$  is the normalized theoretical OCT signal (ranging from 0 to 1). For the convenience of expression, we denote  $\mu_t(r_i)$  as  $\mu$ . A nonlinear least squares (NLS) problem is defined as follow,

$$\hat{\mu}_t(r_i) = \arg \min_{\mu \geq 0} F(\mu), \quad (11)$$

where  $\hat{\mu}_t(r_i)$  is the estimated AC at  $r_i$  and  $F(\mu)$  is the objective function,

$$F(\mu) = \|f(\mu)\|^2. \quad (12)$$

$F(\mu)$  is linearly approximated by the first-order Taylor expansion,

$$F(\mu + \Delta\mu) = \|f(\mu + \Delta\mu)\|^2 \approx \|f(\mu)\|^2 + f'(\mu)\Delta\mu + \frac{1}{2}f''(\mu)(\Delta\mu)^2 \quad (13)$$

where  $f'(\mu)$  and  $f''(\mu)$  are the first and second derivatives of  $f(\mu)$ . Thus, solving  $\mu$  by Eq.(11) is transformed into solving the optimal increment  $\Delta\mu$  in each iteration,

$$\begin{cases} \Delta\hat{\mu} = \arg \min_{\Delta\mu} [m(\Delta\mu)] \\ m(\Delta\mu) = \|f(\mu)\|^2 + f'(\mu)\Delta\mu + \frac{1}{2}f''(\mu)(\Delta\mu)^2 + \frac{1}{2}\gamma(\Delta\mu)^2 \end{cases} \quad (14)$$

where  $\gamma$  is the damping factor being used to punish the excessive increment and  $\Delta\hat{\mu}$  is the optimal increment.

As shown in Fig. 5, the iterative steps for solving Eq. (14) by Levenberg-Marquadt (L-M) optimization are detailed as follows.

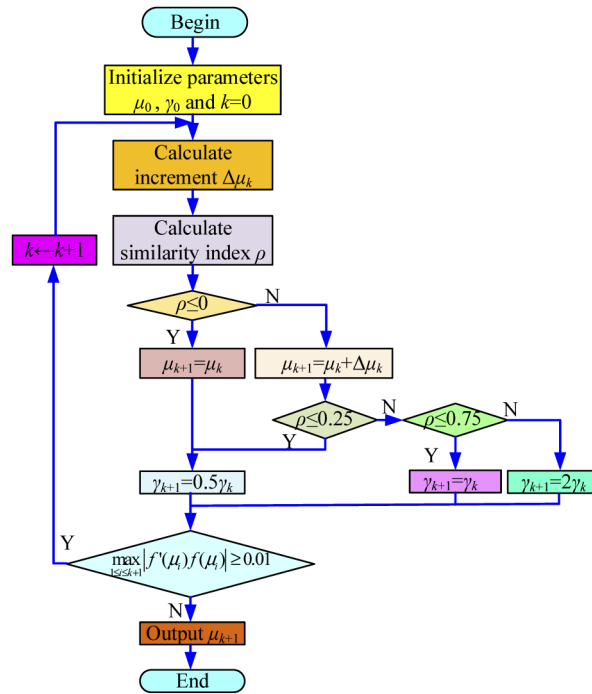


Fig. 5. Flow chart of optimization algorithm.

### Step 1. Initialization.

Initialize the involved parameters including  $\mu_0$ ,  $\gamma_0$ , and  $k=0$ , where  $\mu_0$  is the initial AC at  $r_i$ ,  $\gamma_0$  is the initial damping factor and  $k$  is the number of iteration.  $\mu_0$  can be roughly chosen by referring to the prior histological knowledge about the tissue types.  $\gamma_0$  is usually initialized as 0.01 [35]. In Discussion, we will discuss the dependence of the AC estimates on the initial plans of iteration.

### Step 2. Calculate increment of iteration.

Solve the minimum of  $m(\Delta\mu)$  by letting  $m'(\Delta\mu) = 0$ , thus

$$f'(\mu_k) + f''(\mu_k)\Delta\mu_k + \gamma_k\Delta\mu_k = 0, \quad (15)$$

where  $\mu_k$ ,  $\Delta\mu_k$  and  $\gamma_k$  are the AC, the increment of  $\mu_k$  and the damping factor in the  $k$ -th iteration, respectively. Then,  $\Delta\mu_k$  is determined by

$$\Delta\mu_k = -[f''(\mu_k) + \gamma_k]^{-1}f'(\mu_k). \quad (16)$$

### Step 3. Calculate the similarity index.

Consider

$$F(\mu_k + \Delta\mu_k) \approx \|f(\mu_k)\|^2 + f'(\mu_k)\Delta\mu_k + \frac{1}{2}f''(\mu_k)(\Delta\mu_k)^2 + \frac{1}{2}\gamma_k(\Delta\mu_k)^2 \quad (17)$$

and

$$F(\mu_k) \approx \|f(\mu_k)\|^2, \quad (18)$$

thus

$$F(\mu_k + \Delta\mu_k) - F(\mu_k) \approx f'(\mu_k)\Delta\mu_k + \frac{1}{2}f''(\mu_k)(\Delta\mu_k)^2 + \frac{1}{2}\gamma_k(\Delta\mu_k)^2. \quad (19)$$

Calculate the similarity index that measures the similarity between the linear approximation and the real objective function as follow,

$$\rho = \frac{F(\mu_k + \Delta\mu_k) - F(\mu_k)}{f'(\mu_k)\Delta\mu_k + \frac{1}{2}f''(\mu_k)(\Delta\mu_k)^2 + \frac{1}{2}\gamma_k(\Delta\mu_k)^2}. \quad (20)$$

**Step 4. Adjust the damping factor depending on  $\rho$  and determine  $\mu_{k+1}$ .**

(i)  $\rho \leq 0$  corresponds to a poor linear approximation, thus determine  $\mu_{k+1}$  by

$$\mu_{k+1} = \mu_k. \quad (21)$$

Then, reduce the damping factor, i.e.,  $\gamma_{k+1} = 0.5\gamma_k$ , and proceed to Step 5.

(ii)  $\rho > 0$  corresponds to a good linear approximation, thus determine  $\mu_{k+1}$  by

$$\mu_{k+1} = \mu_k + \Delta\mu_k. \quad (22)$$

Then, adjust the damping factor as follows,

$$\gamma_{k+1} = \begin{cases} 0.5\gamma_k, & 0 < \rho \leq 0.25 \\ \gamma_k, & 0.25 < \rho \leq 0.75 \\ 2\gamma_k, & \rho > 0.75 \end{cases} \quad (23)$$

and proceed to Step 5.

**Step 5. Ending of the iteration.**

If

$$\max_{1 \leq i \leq k+1} |f'(\mu_i)f(\mu_i)| \geq 0.01, \quad (24)$$

proceed to Step 6; otherwise, the iteration ends and output  $\mu_{k+1}$ .

**Step 6. Update  $k$  by  $k \leftarrow k+1$  and turn to Step 2.**

The AC at each pixel position in an image is estimated by above iterative steps. After normalization and pseudocolor coding, the AC matrix is transformed into a parametric image.

### 2.3. Performance evaluation

We evaluated the proposed method experimentally using simulated images and clinical images. The experiments of numerical simulation aim to quantitatively evaluate the accuracy of AC extraction, discuss the influence of the initial plans of iteration and compare with other methods. The experiments on clinical pullback image data aim to demonstrate the validity of our method in retrieving ACs of arterial tissues and characterizing tissue types.

We implemented the method by programming with MATLAB (2016a, The MathWorks, Inc., Natick, MA) on a personal computer configured with a 2.6 GHz Intel Core I7-4510U as CPU, 8GB RAM and Windows 10 64 bits as the operating system. The input data involved in the method are the total number of photons launched, the initial AC, and the central wavelength and FWHM of the incident laser which are provided by the imaging probe.



### 2.3.1. Clinical image preparation

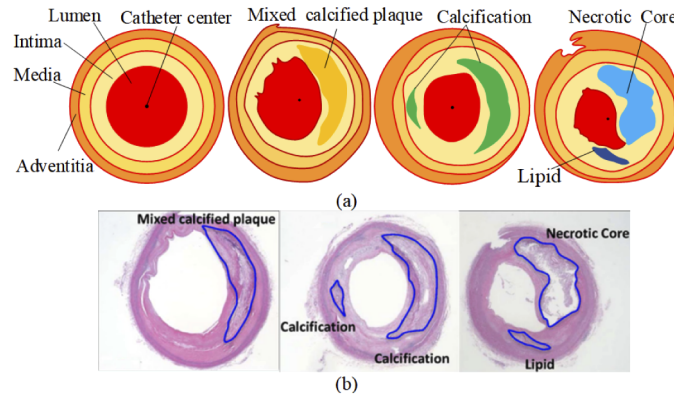
The clinical images were acquired using a commercially available spectral-domain OCT system (C7XR OCT system, LightLab Imaging/St. Jude Medical Inc, St Paul, MN). A 2.7F C7 Dragonfly OCT catheter was inserted into the target vascular lumen and pushed to the distal end through the culprit lesion. It was then drawn back automatically at 20 mm/s, while blood was displaced by an injection of contrast media at rates of 2 to 4 ml/s. The transversal images at an A-line rate of 120 kHz with the axial spatial resolution of 15  $\mu\text{m}$  and transversal resolution of 25  $\mu\text{m}$  were digitally stored for off-line analysis. Table 2 lists the operating and characteristic parameters of the setup.

**Table 2. Parameters of Imaging System in Forward Simulation**

Incident light power (mW)	Noise power (mW)	Central wavelength (nm)	FWHM (nm)	Spectrometer resolution (nm)	Wavelength spacing (nm/pixel)	Maximal detection depth (mm)
2000	$4 \times 10^{-3}$	1310	40	0.11	0.08	2

### 2.3.2. Simulated image preparation

In simulation experiments, we constructed the cross-sectional geometry of computer-simulated arterial phantoms by referring to the histological structures of healthy coronary arterial wall and different types of atherosclerotic plaques. We employed Microsoft Visio (v.2016, Microsoft Corporation, Redmond, State of Washington) to draw the vascular cross sections. Figure 6(a) provides four examples, in which the structures of the arterial wall and plaques are referred to the histology in Fig. 6(b) [18]. The radius of the vascular cross sections is 2.0 mm.



**Fig. 6.** Computer-simulated coronary arterial phantoms. (a) Cross-sectional geometry of the phantoms which are numbered 1#, 2#, 3# and 4# from left to right; (b) Histological cross sections of coronary arterial vessels and atherosclerotic plaques (Hematoxylin and eosin stain) (Ref. [18]).

We simulated IVOCT images basing on the principle of OCT signal formation as described in Sec. 2.1. The instrumental parameters of the OCT system were in accordance with the characteristics of the clinical system described above. We added the Gaussian random noise with the SNR of 55 dB to the simulated OCT signal to be more close to the real situation.

### 2.3.3. Figures of Merit

We adopted the error percentage, energy error depth (EED) [19] and structural similarity (SSIM) [36] to quantitatively evaluate AC estimates. The error percentage is defined by

$$E = (|\hat{\mu}_t - \mu_t| / \mu_t) \cdot 100\%, \quad (25)$$

where  $\mu_t$  and  $\hat{\mu}_t$  are the preset and estimated AC value, respectively. The definition of the error-energy is [19]

$$e(\theta_i) = (||\hat{\mu}_{it} - \mu_{it\text{true}}||_2 / ||\mu_{it\text{true}}||_2) \cdot 100\%, \quad (26)$$

where  $\theta_i \in [0, 2\pi]$ ,  $e(\theta_i)$  is the error-energy in the direction of  $\theta_i$  in the vascular cross section and  $\hat{\mu}_{it}$  and  $\mu_{it\text{true}}$  are the row vectors with length  $M$  which are composed of the estimated and true AC in the direction of  $\theta_i$ , respectively. The EED is defined as the average detection depth when the error-energy exceeds 5% [19]. A larger EED indicates higher estimation accuracy. The SSIM measures the similarity between the images representing the maps of the estimated and true ACs. It is defined by [36]

$$SSIM(X, Y) = \frac{(2\mu_X\mu_Y + c_1)(2\sigma_{XY} + c_2)}{(\mu_X^2 + \mu_Y^2 + c_1)(\sigma_X^2 + \sigma_Y^2 + c_2)}, \quad (27)$$

where  $X$  and  $Y$  are the true and recovered image,  $\mu_X$  and  $\mu_Y$  are the average intensities of  $X$  and  $Y$ ,  $\sigma_X$  and  $\sigma_Y$  are the standard deviation of intensities of  $X$  and  $Y$ ,  $\sigma_{XY}$  is the covariance of intensities of  $X$  and  $Y$ , and  $c_1$  and  $c_2$  are extremely small positive constants being used to ensure the computational stability. A larger SSIM suggests better recovery or estimate.

### 3. Results

#### 3.1. Results of simulated images

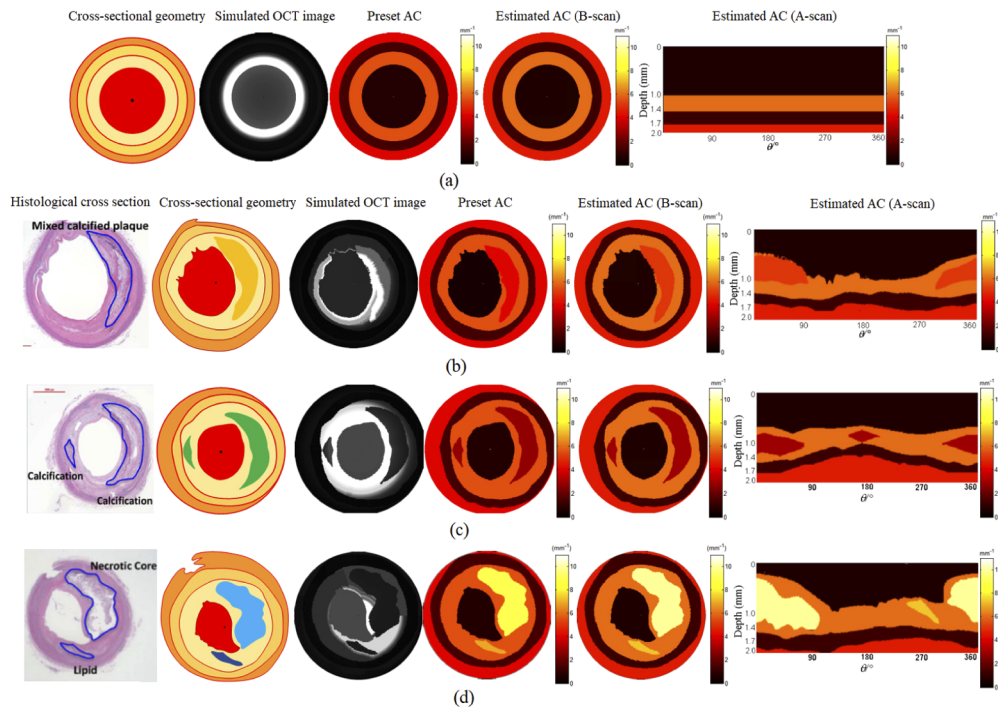
Figure 7 provides the results of the computer-simulated vessel phantoms. The simulated IVOCT images are  $300 \times 300$  pixels (8 mm $\times$ 8 mm) with 8-bit gray levels. For simulation of an A-scan,  $1 \times 10^7$  photons were employed. The A-scans from each 360° rotation were arranged in a B-scan. In the iteration  $\mu_0 = 0.08 \text{ mm}^{-1}$  and  $\gamma_0 = 0.01$ . These results reveal that the AC maps display the multi-layered structures of the vessel wall and plaques with significantly improved contrast in comparison with the structural images. Table 3 lists the true and estimated AC values as well as the error percentages lower than 10%. These convincing results indicate that our method is able to estimate the ACs from IVOCT images with a high accuracy.

**Table 3. Estimated and Preset AC Values and Error Percentages for Simulated Arterial Phantoms**

Tissue type	Radial depth (mm)	AC (mm <sup>-1</sup> )		Error percentage %
		True value	Estimate	
Lumen	0~1.0	0.5	0.452	9.600
Intima	1.0~1.4	4	4.359	8.975
Necrotic core	1.0~1.4	10	10.956	9.562
Lipid	1.0~1.4	5	5.324	6.472
Mixed calcified plaque	1.0~1.4	3	3.226	7.533
Calcification	1.0~1.4	2	2.159	7.950
Media	1.4~1.7	1	0.905	9.500
Adventitia	1.7~2.0	3	3.296	9.867

#### 3.2. Results of clinical images

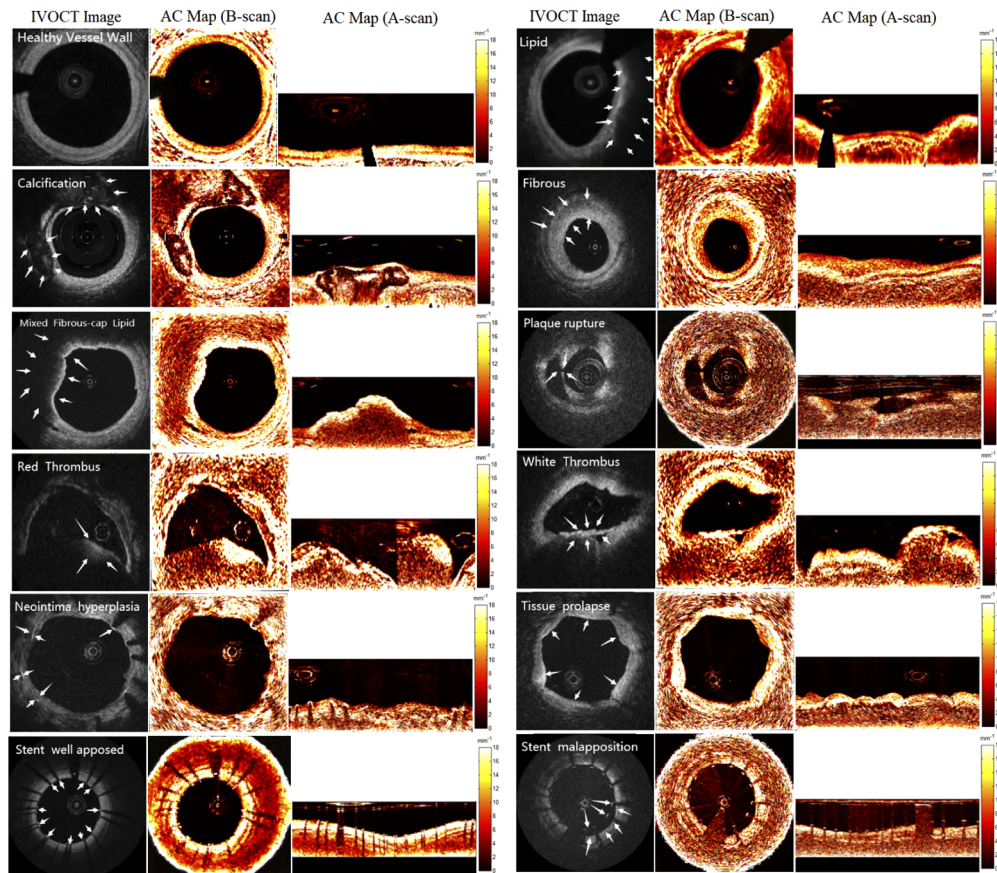
Figure 8 presents the AC maps retrieved from twelve typical frames of clinical IVOCT scanning images ( $300 \times 300$  pixels) containing healthy vessel wall, lesioned tissues and stent struts. The parameter settings were the same as the simulation experiments. The regions of interest (ROIs) including plaque areas and stent struts were manually labeled by experienced physicians. The images of AC mapping provide information regarding the optical properties of different arterial tissue types, which can be applied in the characterization and differentiation of vulnerable



**Fig. 7.** Results of simulated demonstrations. (a) Phantom 1#; (b) Phantom 2# ; (c) Phantom 3#; (d) Phantom 4#.

plaques as well as the guidance and assessment of minimally invasive coronary interventions. These results demonstrate the feasibility of using our method to quantitatively characterize and differentiate at least seven different arterial tissue types, i.e., healthy vessel wall (trilaminar structure of intima (signal-rich or high backscattering), media (signal-poor or low backscattering) and adventitia (heterogeneous and high backscattering)), lipid-rich plaque (low-intensity and irregular signal with an indistinct and irregular border, strong attenuation), calcification (signal-poor or heterogeneous region with a sharp defined border), fibrous plaque (signal-rich or high backscattering, homogeneous and weak attenuation), fibrous cap (a signal-rich tissue layer, overlying a signal-poor region), plaque rupture (fibrous-cap disruption along with cavity formation), thrombus (an irregular mass attached to the luminal surface, including red thrombus (red blood cell-rich, high backscattering and high attenuation) and white thrombus (platelet-rich, low backscattering, homogeneous and low attenuation)) [37–40]. The quantitative tissue characterization by AC extraction is helpful to the qualitative tissue classification that relies on the interpretation of image textures and structural features. Besides, the stents after implantation can be assessed from the AC maps, including neointima hyperplasia, tissue prolapse (tissue protrusion into the lumen between stent struts without apparent surface disruption), and apposition of stent struts (malapposition: the axial distance between the strut's surface to the luminal surface is greater than the strut thickness; well apposed: the distance is less than the strut thickness [37]).

Table 4 presents the ranges of the histological Refs. [7,39] and estimated AC values of different tissue types. In the table, the ranges of the ACs were measured from the manually selected regions in the AC maps shown in Fig. 8. These results indicate that the AC estimates obtained by our method are basically consistent with the histological references.



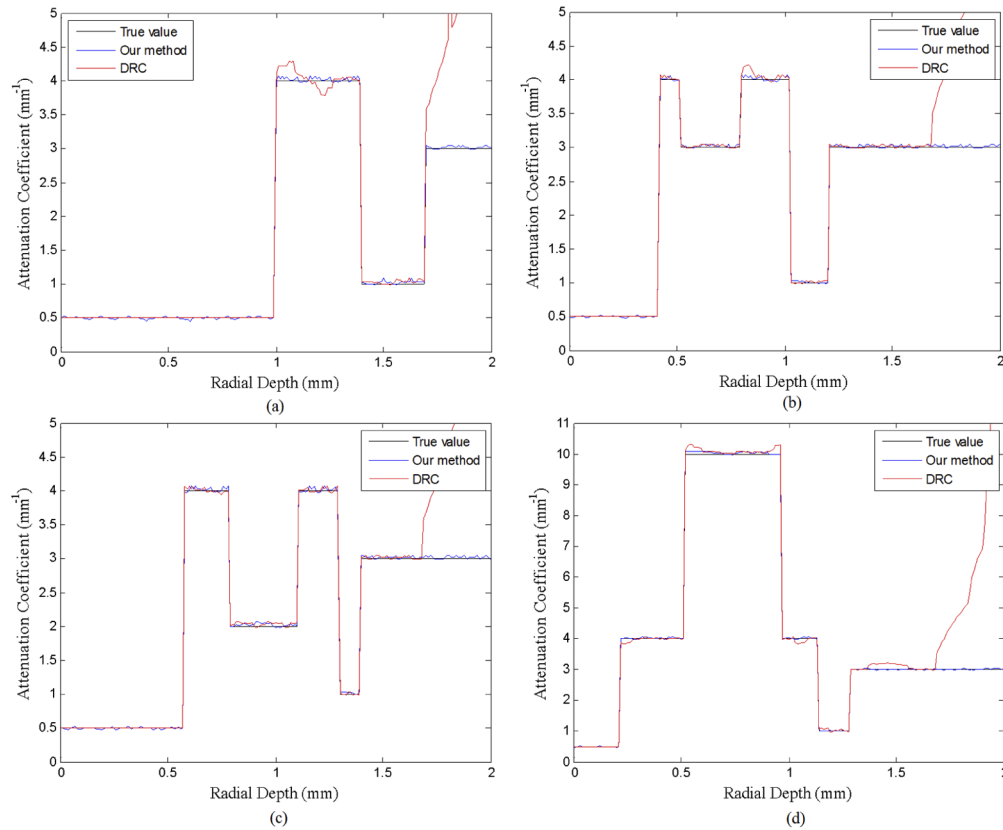
**Fig. 8.** Clinical IVOCT images and images of AC along with manually labeled ROIs.

**Table 4.** Comparison of Estimates of AC from Clinical Images with Histological References (Ref. [7,39])

Tissue name	AC ( $\text{mm}^{-1}$ )	
	Histological reference	Estimate
Healthy arterial wall	2~5	2.392~4.454
Calcified plaque	2~7.1	1.894~7.517
Fibrous plaque	1.07~7.6	1.611~6.214
Lipid plaque	7.62~18.2	5.127~17.162
White thrombus	5.2~17.5	3.120~15.652
Red thrombus	3.6~9.6	5.852~8.924

### 3.3. Comparison results

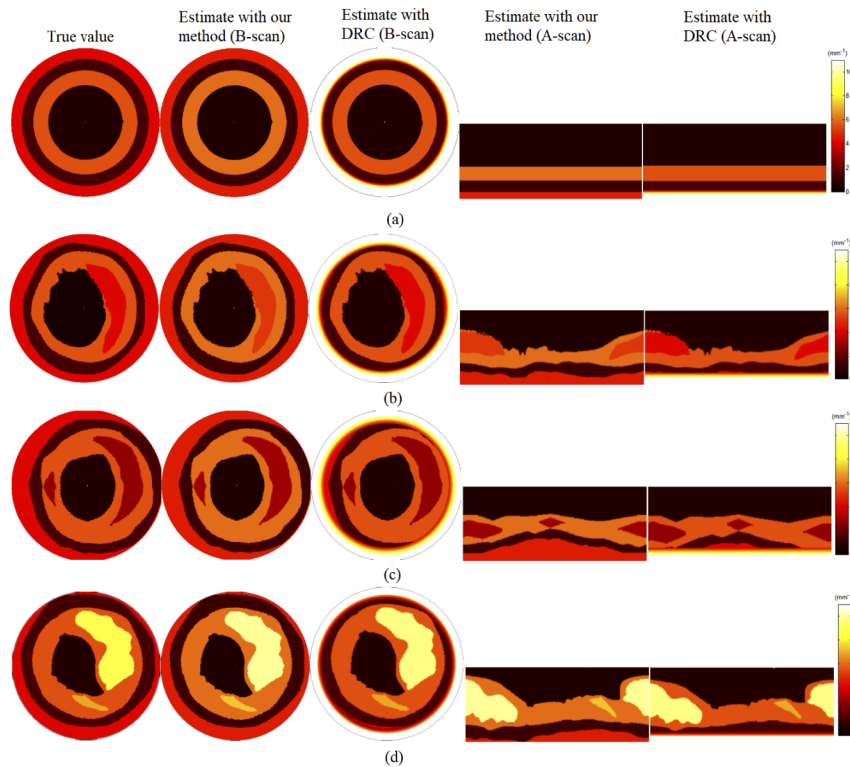
In this section, the results of comparing our method with Smith's DRC method [19] are given. In our method, the initial plan of iteration is  $\mu_0=0.08 \text{ mm}^{-1}$  and  $\gamma_0=0.01$ . DRC utilizes intensity weighted horizontal total variation (iwhTV) regularization to optimize the estimated ACs, where the regularization parameter is  $\eta=3$ . Figures 9–11 present the results of the simulated images and Fig. 12 provides the results of the clinical images.



**Fig. 9.** AC values along one radius of the vessel phantoms estimated by using our method and DRC. (a) Phantom 1# (90°); (b) Phantom 2# (105°); (c) Phantom 3# (190°); (d) Phantom 4# (75°).

In our simulated phantom preparation, not all the incident light is assumed to be completely attenuated at depths ranging from 0 to 2.0 mm. As indicated in Fig. 9, the overestimation at the end of the depth profile (i.e., at depths ranging from 1.7 mm to 2.0 mm) is observed in the estimates by DRC because of its premise of full attenuation. This conclusion can also be drawn from the visualization of the AC maps shown in Fig. 10. In contrast, the estimates at the end of the depth profile by using our method are closer to the true values than the DRC estimates. Figure 11 suggests that both the EED and SSIM of the estimates by our method are improved by up to 30% and 10% in comparison with those by DRC. The effective range of the AC estimation with our method is the entire radial depth, i.e., 2.0 mm, while that of DRC is about 1.7 mm. Our method has a relatively higher time cost than DRC because of its numerical iterative nature in contrast to DRC's analytical nature. Fortunately, this computational burden can be greatly reduced by using parallel codes and a high-performance parallel computing platform that exploits graphics processing units (GPUs).

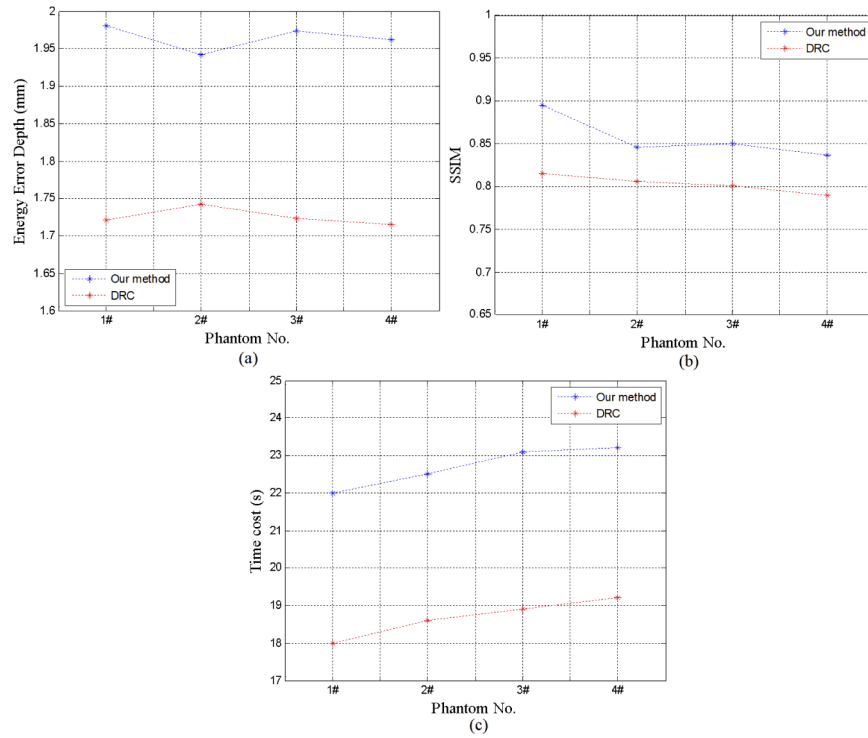




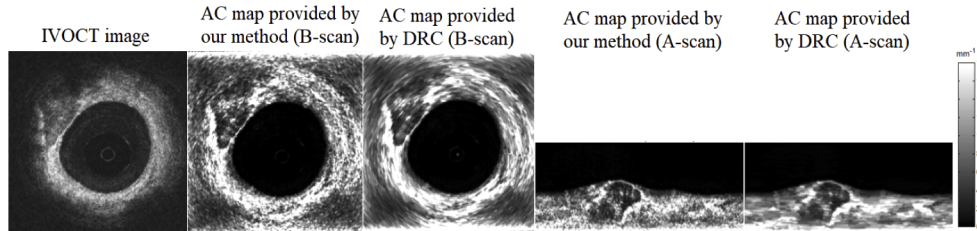
**Fig. 10.** AC maps of the true value and estimates by our method and DRC. (a) Phantom 1#; (b) Phantom 2#; (c) Phantom 3#; (d) Phantom 4#.

Moreover, DRC has been demonstrated to be efficient in extracting ACs in horizontal layered tissues such as bladder, skin and retina by utilizing the correlation of A-scan data to smooth tissue layers [19]. However, coronary arteries with atherosclerotic lesions usually present complicated anatomical structures and have various tissue compositions as well. Hence, a vascular cross section in a transverse structural image may involve multiple tissue types, showing complex structures instead of horizontal layered structure. Results in Fig. 12 are consistent with this conclusion, where the transverse ambiguity is observed in the DRC estimates. Our method outperforms DRC in the contrast of different tissues in the AC maps, especially at the end of the depth profile.





**Fig. 11.** Comparison of figures of merit and time cost of AC estimates by our method and DRC. (a) EED; (b) SSIM; (c) Time cost.

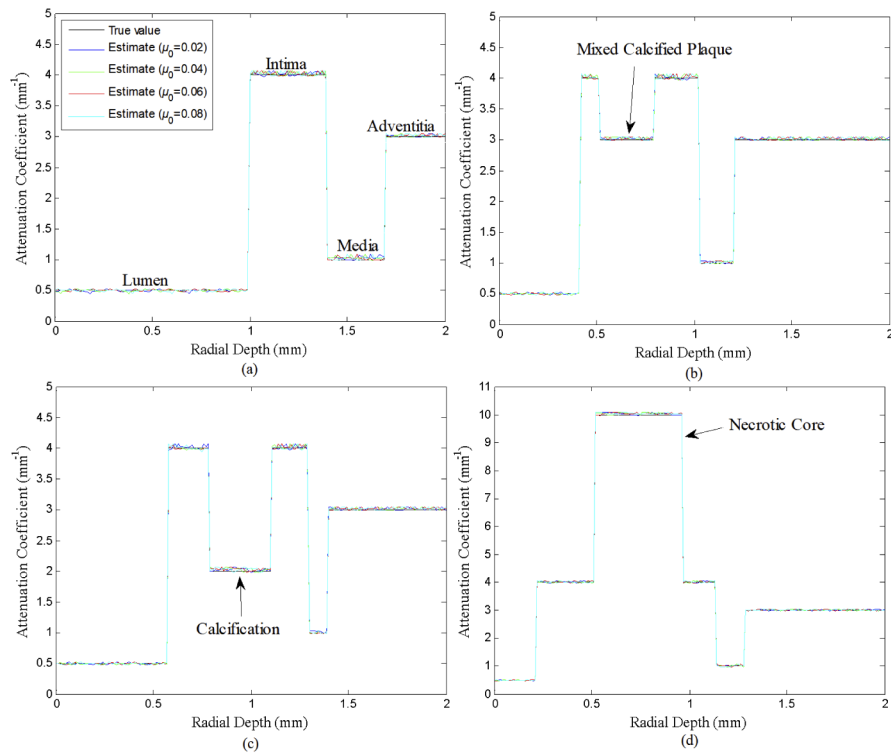


**Fig. 12.** Clinical IVOCT images and AC maps provided by our method and DRC.

## 4. Discussion

### 4.1. Dependence on initial plans of iteration

We discussed the influence of the initial plans of iteration on the AC extraction based on the simulation results. First, we compared the AC estimates obtained by changing the initial AC  $\mu_0$ , while the damping factor remained unchanged, i.e.,  $\gamma_0=0.01$ . Figure 13 suggests small differences in the AC estimates when  $\mu_0=0.02 \text{ mm}^{-1}$ ,  $0.04 \text{ mm}^{-1}$ ,  $0.06 \text{ mm}^{-1}$  and  $0.08 \text{ mm}^{-1}$ . Figure 14 indicates that the EED and SSIM of the AC estimates when  $\mu_0=0.08 \text{ mm}^{-1}$  are slightly higher than those of other cases with the least times of iteration. Second, we set the damping factor  $\gamma_0=0.01, 0.02, 0.03$  and  $0.04$  respectively, while  $\mu_0=0.08 \text{ mm}^{-1}$ . The results in Fig. 15 reveal that the lowest iteration times and time cost are achieved when  $\gamma_0=0.01$ . In the L-M optimization, the damping factor is adjusted dynamically during iteration, so it has major influence on the iteration times and time cost, but minor influence on the iteration result.



**Fig. 13.** Estimated values of AC in the case of different  $\mu_0$ . (a) Phantom 1# (90°); (b) Phantom 2# (75°); (c) Phantom 3# (105°); (d) Phantom 4# (190°).

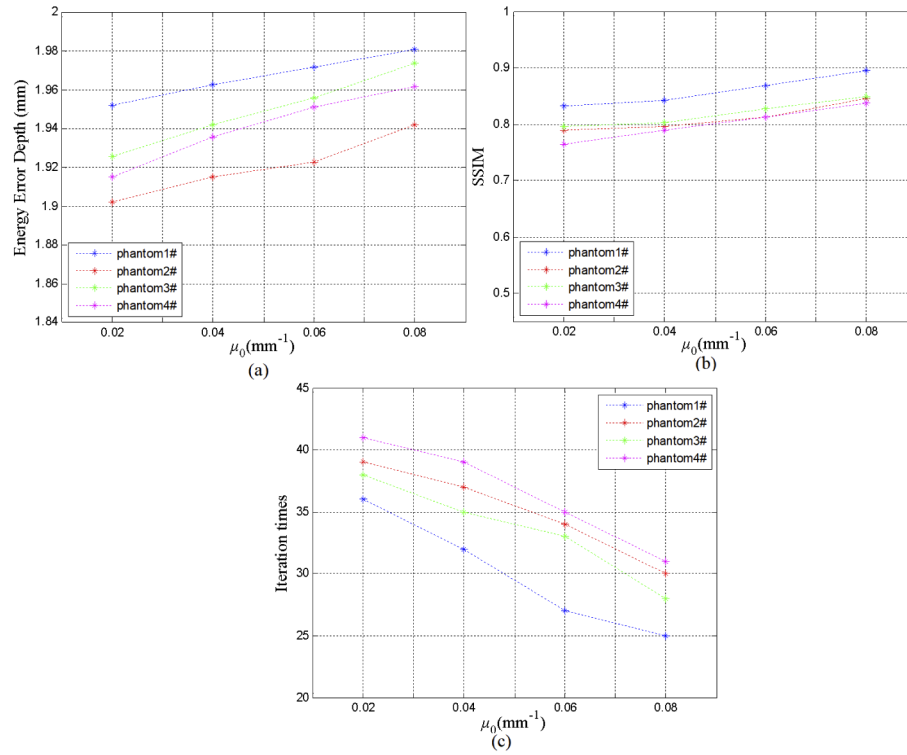
## 4.2. Discussion on model-based scheme

The AC estimation from OCT measurements is a distributed parameter estimation problem. It is essentially the inversion of a light transport operator. We designed a model-based minimization scheme for the recovery of AC as a function of position from IVOCT images. Two main aspects are involved in this scheme: forward model and optimization framework. We will discuss both issues in this section.

### 4.2.1. Model selection

The selection of the OCT signal model is critical in extracting AC from OCT data and it should be based on the chemical compositions and biological characteristics of tissue regions of interest. Considering the simplicity of the SS model, many efforts have been made to improve the accuracy and adaptability of the SS-based schemes such as CF and DR methods. For IVOCT, clinical applications have demonstrated that the effects of multiple scattering should be considered in coronary arterial vessel walls with multi-layered structures and atherosclerotic plaques [3,29]. Hence, AC measurements with higher accuracy and stability will be obtained by taking into account multiple scattering of photons in comparison with single scattering although it may be computationally intensive.

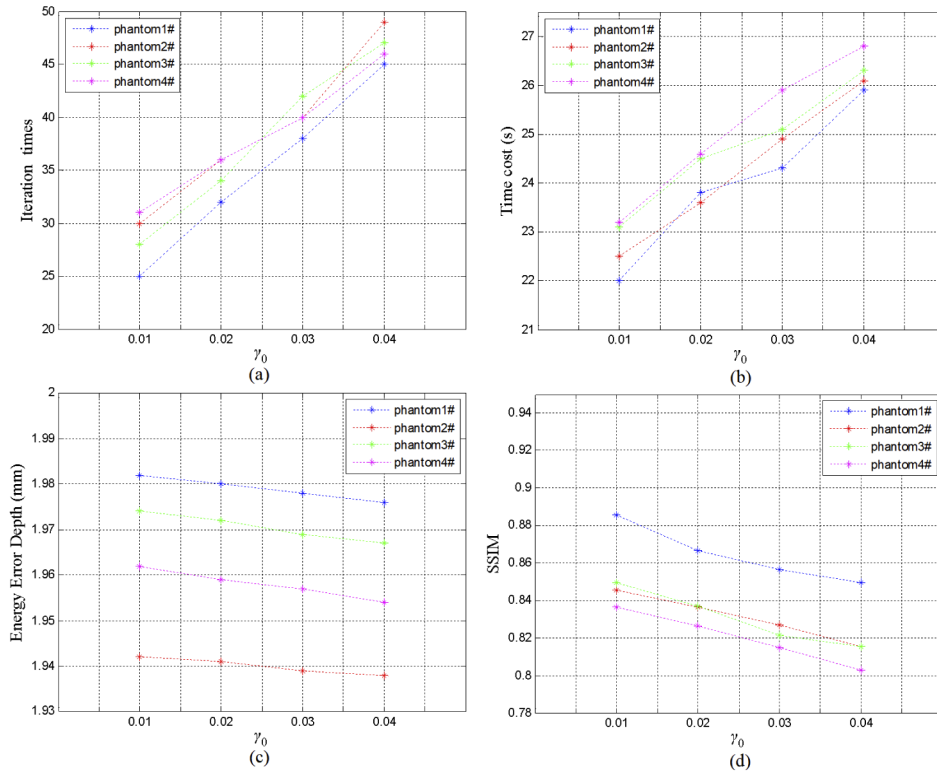
When light propagating in a turbid medium, the attenuation is mainly caused by absorption and scattering which are parameterized by the absorption coefficient  $\mu_a$  and scattering coefficient  $\mu_s$  of tissues, respectively. The sum of  $\mu_a$  and  $\mu_s$ , i.e., the AC  $\mu_t$ , represents tissue optical properties. In OCT operating in a NIR spectrum window, the attenuation is dominated by the scattering while the absorption by those optical absorbers is very low to achieve high enough imaging depths.



**Fig. 14.** Figures of merit of AC maps obtained in the case of different  $\mu_0$ . (a) EED; (b) SSIM; (c) Iteration times.

In the CF- and DR-based methods, the loss of the backscattered signal with the penetration depth caused by optical absorption and scattering is parameterized by a decay coefficient  $\mu_{\text{OCT}}$  and then is modeled as a single exponential decay governed by the Lambert-Beer law. This simple parameterization greatly alleviates the complexity, but sacrifices the direct relationship with the optical properties of tissues (e.g., anisotropy factor  $g$ ,  $\mu_a$  and  $\mu_s$ ) to achieve robust and unique measurement [4]. Moreover, the estimated  $\mu_{\text{OCT}}$  may be theoretically lower than  $\mu_t$  in the presence of multiple scattering events because more light will be detected than expected by the SS model due to multiple scattering.

In the model-based minimization scheme of a parameter estimation problem, the forward signal model is required to be sufficiently accurate to capture the essential characteristics of the light field. Besides, it should be fast enough computationally to make its use in iterative inversion methods, where it may need to be evaluated many times [41]. In biomedical optics, the MC method represents the gold standard for modeling light transport in biological tissues, both due to its accuracy and its flexibility in modeling realistic, heterogeneous, multilayered turbid media [42]. Since both the effects of multiple scattering and dependent scattering play an important role in the quantification of optical tissue properties [43], we will take into account the dependent scattering by utilizing the well-established MC simulation software in our future work as has been done in [42]. Moreover, because of the inherently parallel nature of MC simulations, the time cost could be reduced significantly by implementing the simulator on GPUs using the compute unified device architecture (CUDA) programming language [44] as the availability of large computing clusters and clusters of GPUs is increasing quickly [31].



**Fig. 15.** Figures of merit of AC maps obtained in the case of different  $\gamma_0$ . (a) Iteration times; (b) Time cost; (c) EED; (b) SSIM.

#### 4.2.2. Optimization framework

We used the least-squares minimization, a common and well-developed framework to solve the inverse problem, where we solved the NLS problem with a L-M algorithm. The initial plans of iteration are required to be properly selected according to the prior histological knowledge of the tissue types and components. It may lower the automation of the method. In addition, as an iterative scheme, the computational burden that might limit the real-time applications should be considered. An efficient optimization tool such as Bregman [45] and fast iterative shrinkage thresholding algorithm (FISTA) [46], parallel codes and a high-performance parallel computing platform that exploits GPUs being used to reduce the computational burden may improve the efficiency of our method.

In recent years, deep learning (DL) has shown great potential in the field of medical imaging with the rapid growth of high-performance processors and big data technology. DL, especially convolutional neural network (CNN), has been proven an efficient tool in medical image analysis, processing and reconstruction [47]. Many efforts have also been made in DL-based approaches of quantitative parametric imaging, such as quantitative photoacoustic imaging which aims to solve an ill-posed inverse problem of recovering optical absorption coefficient and scattering coefficient from acoustic measurements [48]. Big data is crucial for DL. IVOCT has been routinely applied in minimally invasive cardiac interventions, so that a large amount of clinically acquired image data is available, which can be used to build training data sets. Hence, in the future work, we plan to develop an automated data-driven DL method for quantitative IVOCT by constructing a recovery function in the form of CNN to map an input structural image to an output parametric image of attenuation.

## 5. Conclusion

We proposed a novel method based on error-minimization as an efficient tool of retrieving pixel-wise AC distribution in vascular cross sections from IVOCT images to implement parametric imaging of optical attenuation. We solved the NLS problem regarding the measured and modeled backscattered signal being obtained by forward MC modeling. Results of simulated data demonstrated the accuracy of AC extraction with the error percentage less than 10%. Results of clinical images demonstrated the consistency of AC estimates of different tissue types with the histological references. The parametric imaging of AC enables the differentiation and characterization of at least seven tissue types, i.e., healthy vessel wall, lipid-rich plaque, calcified plaque, fibrous plaque, fibrous cap, plaque rupture and red/white thrombus, as well as stent apposition (well apposition, malapposition, intima hyperplasia and plaque prolapse). Comparison results validate the improvement of up to 30% in the EED and 10% in the SSIM of the parametric imaging of attenuation in comparison with the DRC approach.

One issue that must be pointed out that we evaluated the accuracy of our method quantitatively with computer-simulated vessel phantoms. However, it is evident that an AC estimation scheme should be further evaluated through validation measurements of tissue phantoms with well specified attenuation coefficients as has been done in [49]. It will be another work for us in the further study.

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